

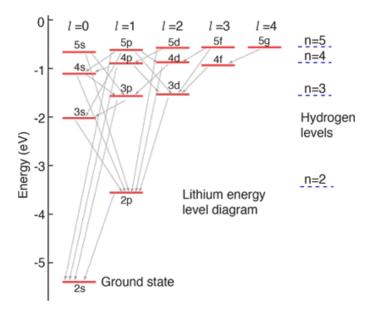


Adapting the ab-initio IMSRG for open-shell atomic systems

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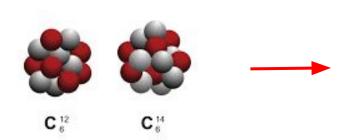
$$H = \sum_{i}^{N} -rac{
abla_{i}^{2}}{2} - rac{Z}{r_{i}} + \sum_{i < j}^{N} rac{1}{r_{ij}}$$



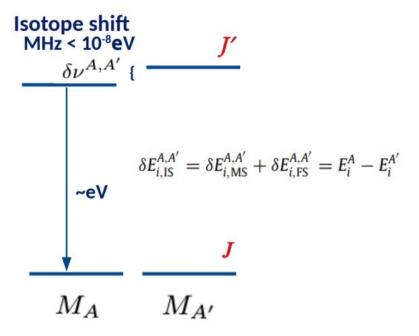
- → Many-body atomic hamiltonian assuming point-like nucleus
- → Get spectrum can be experimentally measured by spectroscopy

$$rac{1}{r}
ightarrow
ho_{nuc}(r)$$

→ Nuclear structure modifies Coulomb potential



→ Differing nuclear structure among isotopes shifts spectra - experimentally measured by laser spectroscopy



$$\delta E_{i,\text{MS}}^{A,A'} = \left(\frac{M' - M}{MM'}\right) (K_{i,\text{NMS}} + K_{i,\text{SMS}}) = \left(\frac{M' - M}{MM'}\right) K_{i,\text{MS}}$$

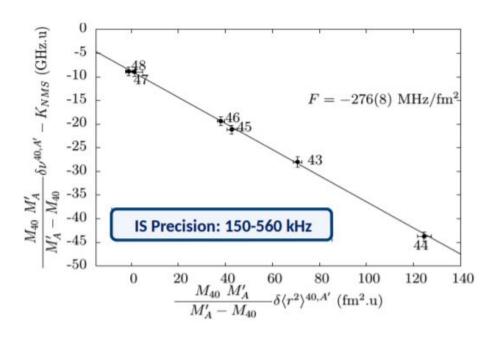
$$\hat{H}_{NMS} = rac{1}{2M} \sum_{j=1}^{N} \left(\mathbf{p}_{j}^{2} - rac{lpha Z}{r_{j}} lpha_{j} \cdot \mathbf{p}_{j}
ight)$$

$$\hat{H}_{NMS} = rac{1}{2M} \sum_{j=1}^N \left(\mathbf{p}_j^2 - rac{lpha Z}{r_j} lpha_j \cdot \mathbf{p}_j
ight) \qquad \qquad \hat{H}_{SMS} = rac{1}{2M} \sum_{j
eq k}^N \left(\mathbf{p}_j \cdot \mathbf{p}_k - rac{lpha Z}{r_j} lpha_j \cdot \mathbf{p}_k
ight)$$

Field Shift:

$$\delta E_{i,\mathrm{FS}}^{(1)A,A'} = -\int_{\mathbf{R}^3} \left[V^A(\mathbf{r}) - V^{A'}(\mathbf{r}) \right]
ho_i^e(\mathbf{r}) d^3\mathbf{r}, \;\; pprox \sum_n F_{i,n} \delta \langle r^{n+2}
angle^{A,A'}$$

$$\delta
u^{A,A'} = K_{MS} \frac{M_{A'} - M_A}{M_{A'} M_A} + F \delta \langle r^2 \rangle^{A,A'}$$



Ca:<r²>:Garcia Ruiz et al. Nature Phys. 12, 594 (2016)

→ Non-linearities in "King plots" point to new electron-nucleon interactions.

Flambaum et al. Phys Rev A 97, 032510 (2018)

→ Laser spectroscopy can be used to probe BSM physics

Stadnik et al. Phys Rev Lett 120, 223202 (2018) - Long range neutrino mediated forces

Similarity Renormalization Group (SRG)

Basic Idea - Diagonalize a matrix by performing continuous Unitary transformations on it, using a suitable generator

$$H(s) = U(s)HU^{\dagger}(s) \equiv H^{\mathrm{d}}(s) + H^{\mathrm{od}}(s) \to H^{\mathrm{d}}(\infty)$$

→ Unitarity preserves inner products - eigenvalues unchanged

$$\frac{\mathrm{d}H(s)}{\mathrm{d}s} = \left[\eta(s), H(s)\right] \qquad \qquad \eta_{I}(s) = \left[H^{\mathrm{d}}(s), H^{\mathrm{od}}(s)\right]$$

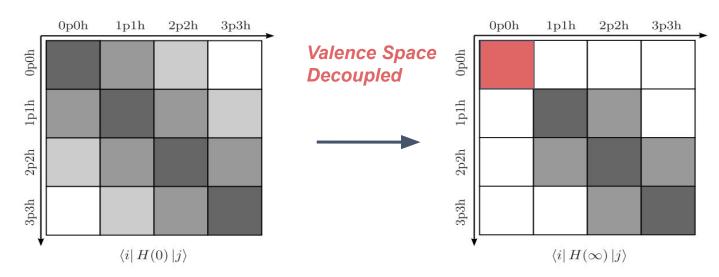
Flow equation for generator - $\eta(s)$

"Wegner Generator"

Valence Space In-Medium SRG

$$H^{\text{od}} = \langle p|H|h\rangle + \langle pp|H|hh\rangle + \dots + \text{h.c.}$$

→ Redefine Hamiltonian using particle-hole excitations from Hartree-Fock ground state.

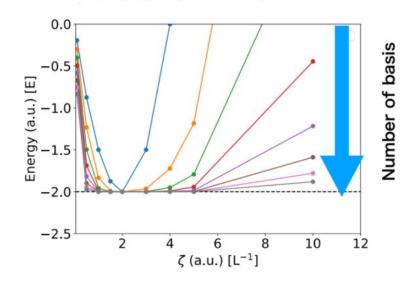


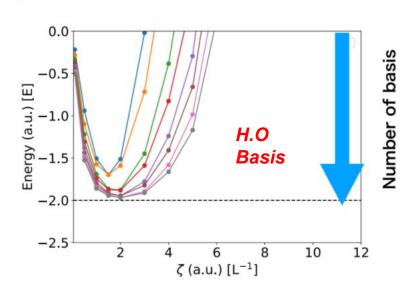
Setup for Atomic Systems

Choice of basis function: *Laguerre function* + Spherical Harmonics

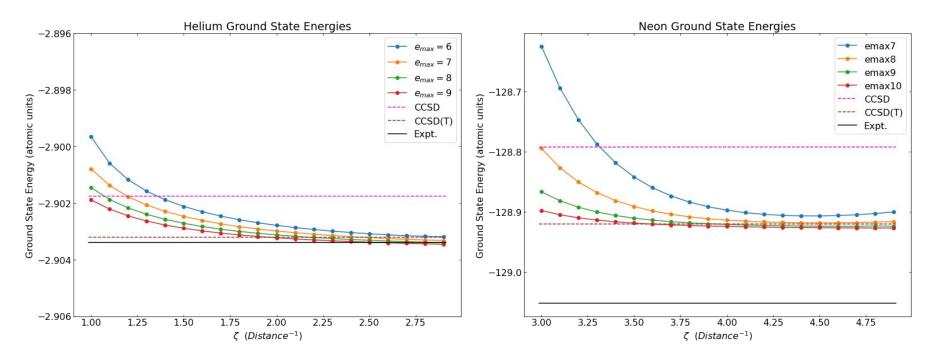
$$R_{nl}(r) = \sqrt{\left(\frac{2\zeta}{a_0}\right)^3 \frac{n!}{(n+2l+2)!}} x^l e^{-x/2} L_n^{2l+2}(x), x = \frac{2\zeta}{a_0} r$$

A. E. McCoy and M. A. Caprio, J. Math. Phys. 57, (2016)



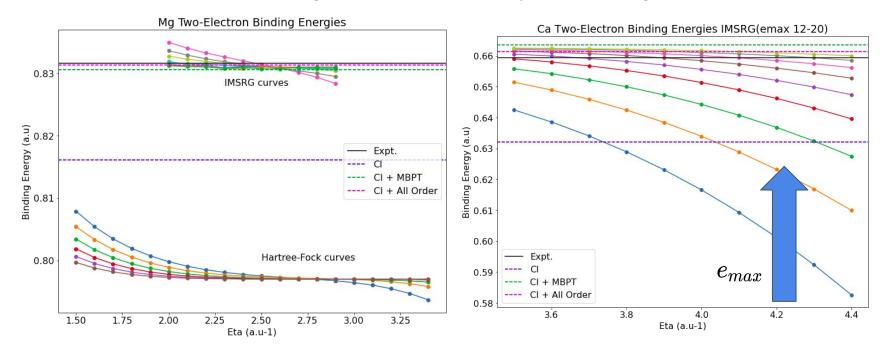


Results - Closed shell



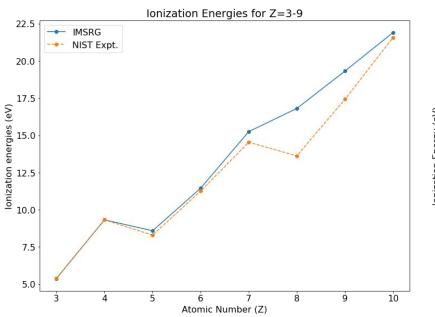
Results - Open shell

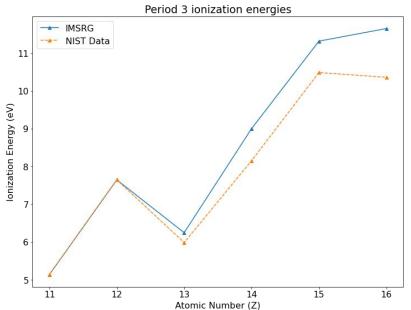
→ Valence space diagonalization done by k-shell to get excited states



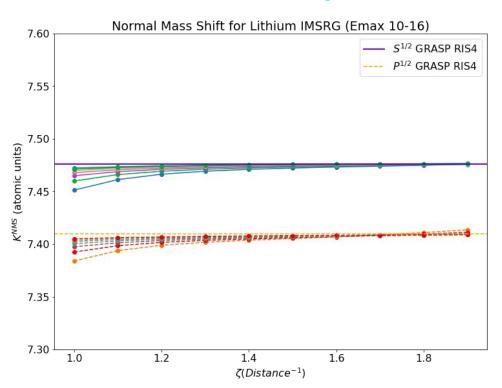
M. S. Safronova et al. Phys. Rev. A 80, 012516 (2009)

Results - Open Shell



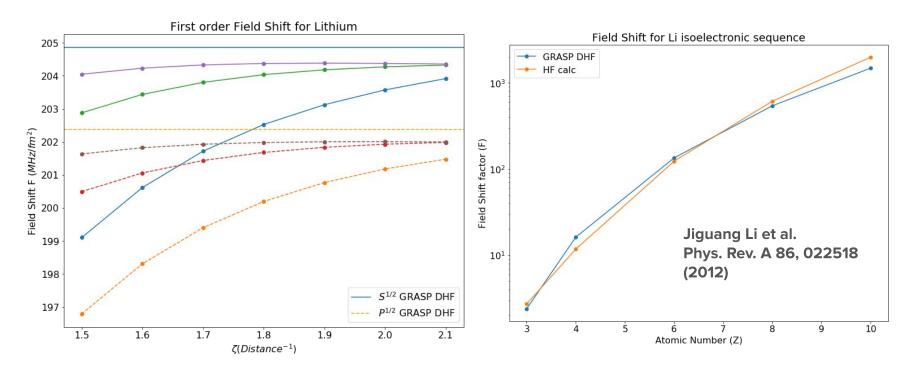


Results - Isotope Shifts



→ Results comparable to current state of the art Relativistic Isotope Shift program (RIS4)
Ekman et al. (2019)

Results - Isotope Shifts



Future Work

→ Many body relativistic corrections - Breit Hamiltonian (Partial progress)

$$H=H_{el}+H_{rel}+H_D+H_{SSC}+H_{OO}+H_{SO}+H_{SSD}$$

- Spin-Orbit / Orbit-Orbit interactions
- → Higher order corrections to isotope shifts
- → Incorporate nuclear moments for hyperfine structure effects.

$$W_{F,J}^{M1} = A_{\rm hf} \mathbf{I}.\mathbf{J}$$

$$W_{F,J}^{E2} = B_{\rm hf} \frac{3(\mathbf{I}.\mathbf{J})^2 + \frac{3}{2}(\mathbf{I}.\mathbf{J}) - I(I+1)J(J+1)}{2I(2I-1)J(2J-1)}.$$

Summary

- → Isotope shifts in atomic spectra allow probing of nuclear structure through laser spectroscopy.
- → VS-IMSRG is a robust ab-initio method to diagonalize atomic many-body Hamiltonians, comparable to Coupled Cluster theory, MBPT.
- → Promising results for spectra, isotope shift factors with potential for further extensions (fine/hyperfine structure).

Thank You! Merci!

IMSRG equation

 $W_{ijklmn} = \langle ijk | \hat{v}^{[3]} | lmn \rangle$.

$$\begin{split} \hat{H} &= E + \sum_{ij} f_{ij} \left\{ a_i^\dagger a_j \right\} + \frac{1}{4} \sum_{ijkl} \Gamma_{ijkl} \left\{ a_i^\dagger a_j^\dagger a_l a_k \right\} + \frac{1}{36} \sum_{ijklmn} W_{ijklmn} \left\{ a_i^\dagger a_j^\dagger a_k^\dagger a_n a_m a_l \right\} \\ &= \left(1 - \frac{1}{A} \right) \sum_a \langle a|\hat{t}^{[1]} |a\rangle n_a + \frac{1}{2} \sum_{ab} \langle ab| \frac{1}{A} \hat{t}^{[2]} + \hat{v}^{[2]} |ab\rangle n_a n_b + \frac{1}{6} \sum_{abc} \langle abc| \hat{v}^{[3]} |abc\rangle n_a n_b n_c \,, \\ &f_{ij} = \left(1 - \frac{1}{A} \right) \langle i|\hat{t}^{[1]} |j\rangle + \sum_a \langle ia| \frac{1}{A} \hat{t}^{[2]} + \hat{v}^{[2]} |ja\rangle n_a + \frac{1}{2} \sum_{ab} \langle iab| \hat{v}^{[3]} |jab\rangle n_a n_b \,, \\ &\Gamma_{ijkl} = \langle ij| \frac{1}{A} \hat{t}^{(2)} + \hat{v}^{[2]} |kl\rangle + \sum_a \langle ija| \hat{v}^{[3]} |kla\rangle n_a \,, \end{split}$$

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